

Propagación y Transmisión Inalámbrica

Dpto. Teoría de la Señal y Comunicaciones
Universidad Carlos III de Madrid

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Fundamentos y conceptos básicos de radiación

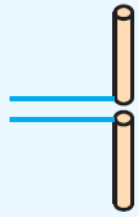
- Introduction to electromagnetic radiation
- The short dipole

The logo for Cartagena99, featuring the text 'Cartagena99' in a stylized font with a blue and orange gradient background.

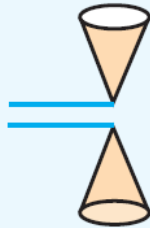
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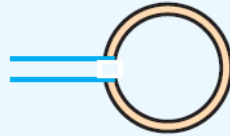
Algunos tipos de antenas



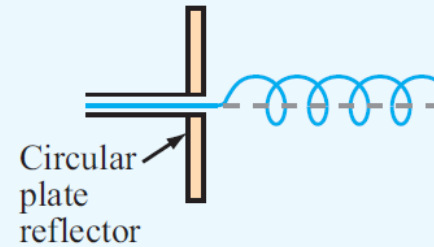
(a) Thin dipole



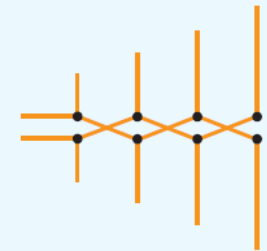
(b) Biconical dipole



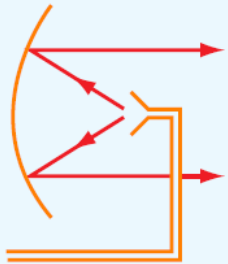
(c) Loop



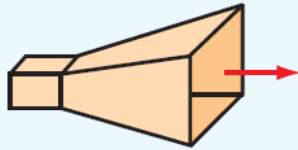
(d) Helix



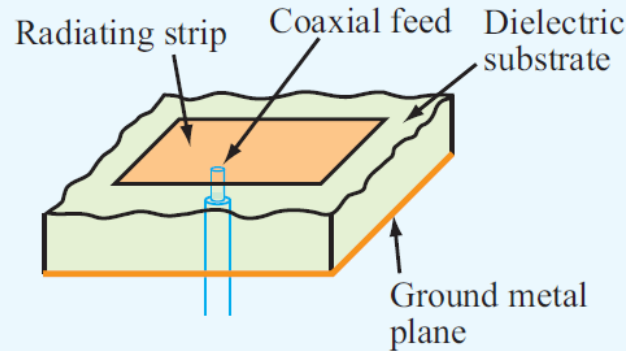
(e) Log-periodic



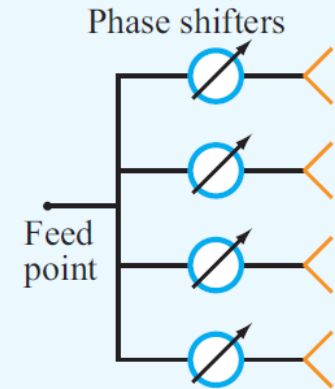
(f) Parabolic dish reflector



(g) Horn



(h) Microstrip

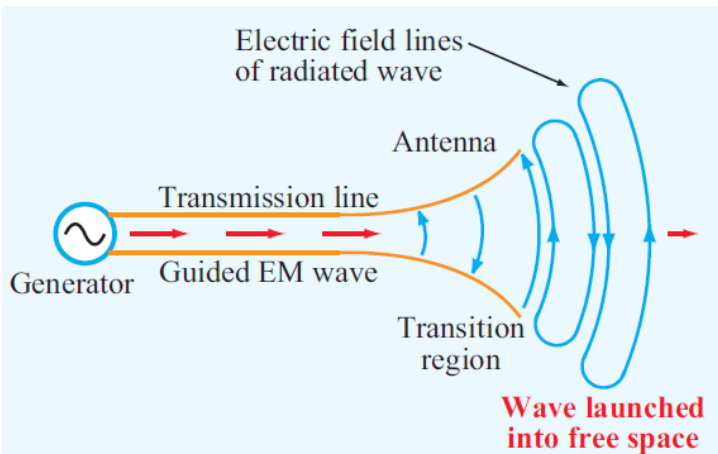


(i) Antenna array

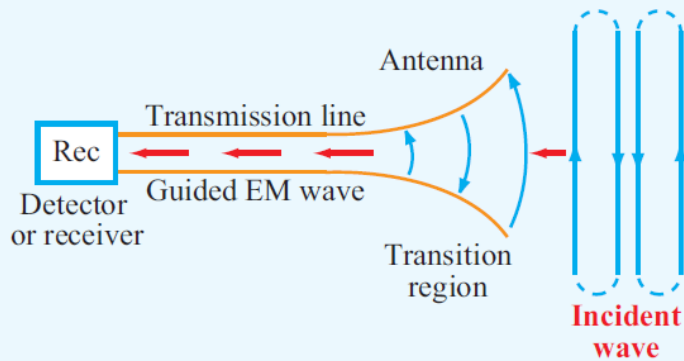
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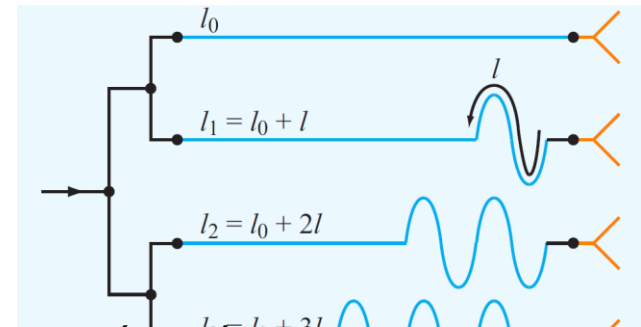
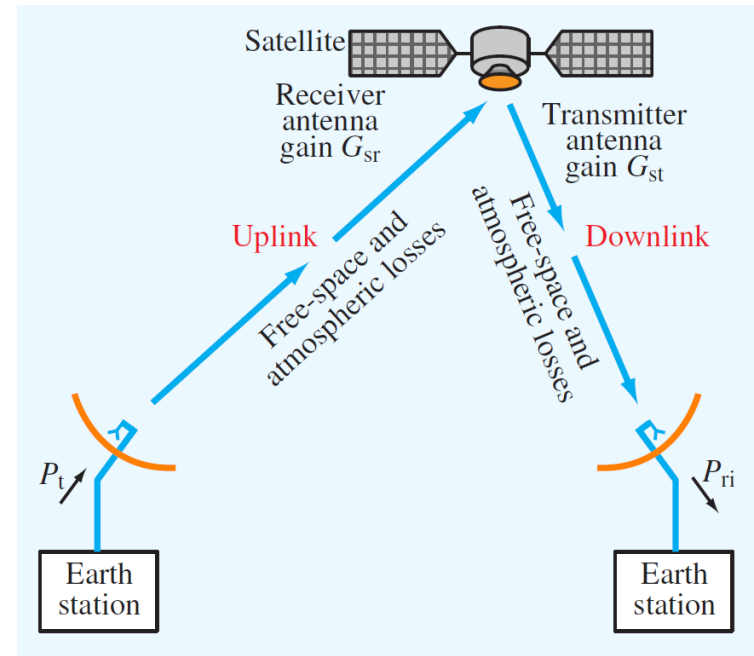
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(a) Transmission mode



(b) Reception mode



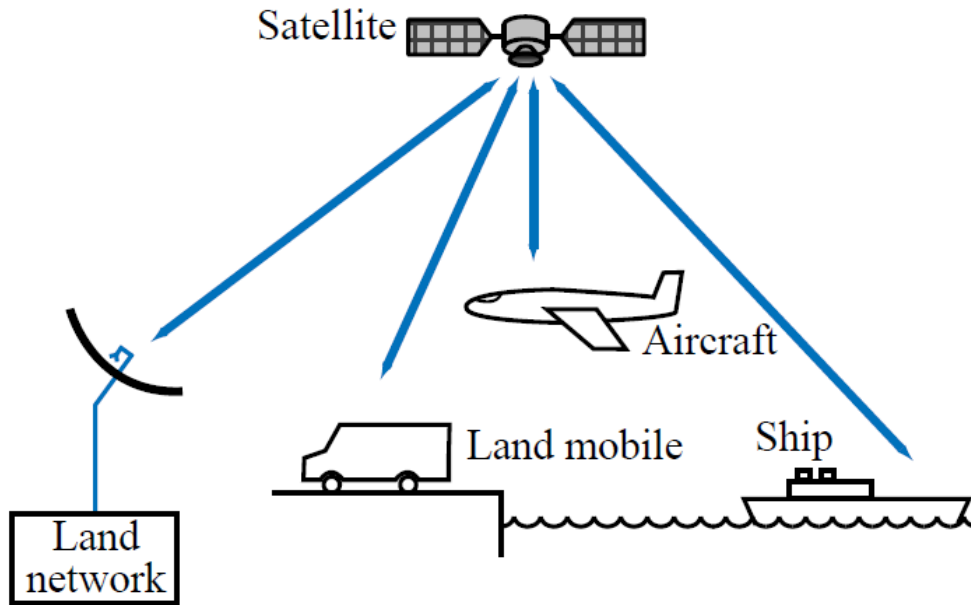
Figuras del libro de Ulaby:
F.T.Ulaby , U.Ravaioli: "Fundamentals of Applied

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Why do we need antennas? Motivation I

- Satellite Communication Systems



http://www.esa.int/spaceinimages/Images/2014/03/Artist_s_impression_of_Sentinel-1A

Vast communication network...voice, data, video services....

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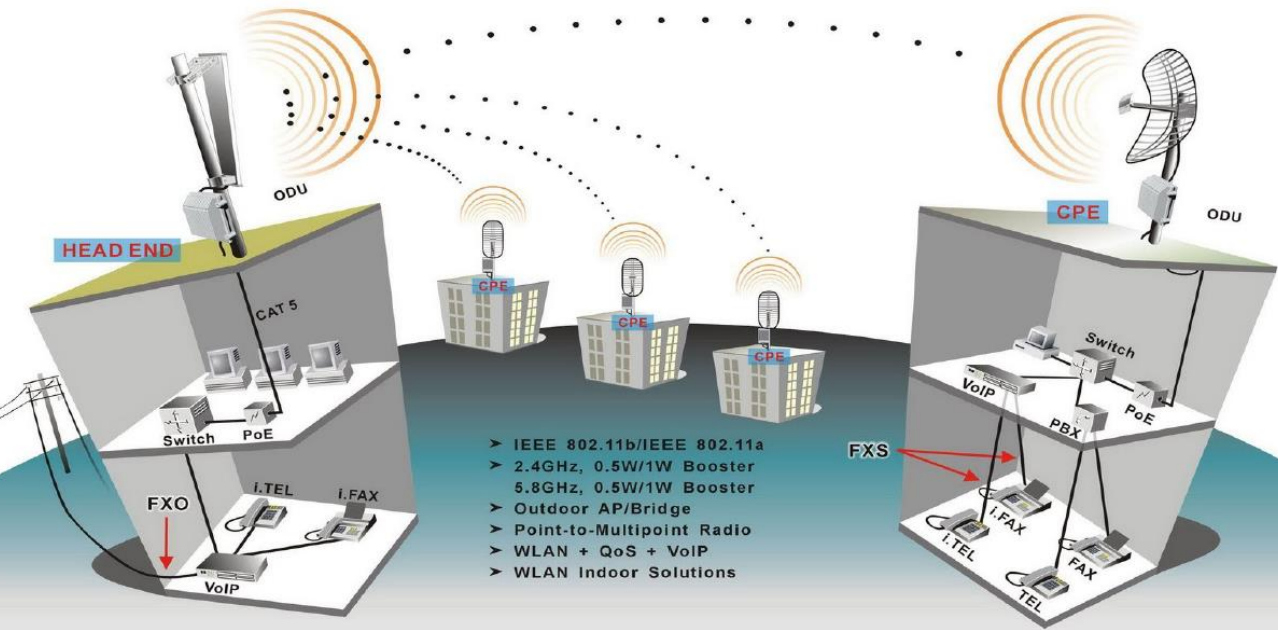
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Why do we need antennas? Motivation II

- Wireless systems....mobile terminals



Antennas are needed and used in wireless systems...

Wireless ground-based telecommunication systems

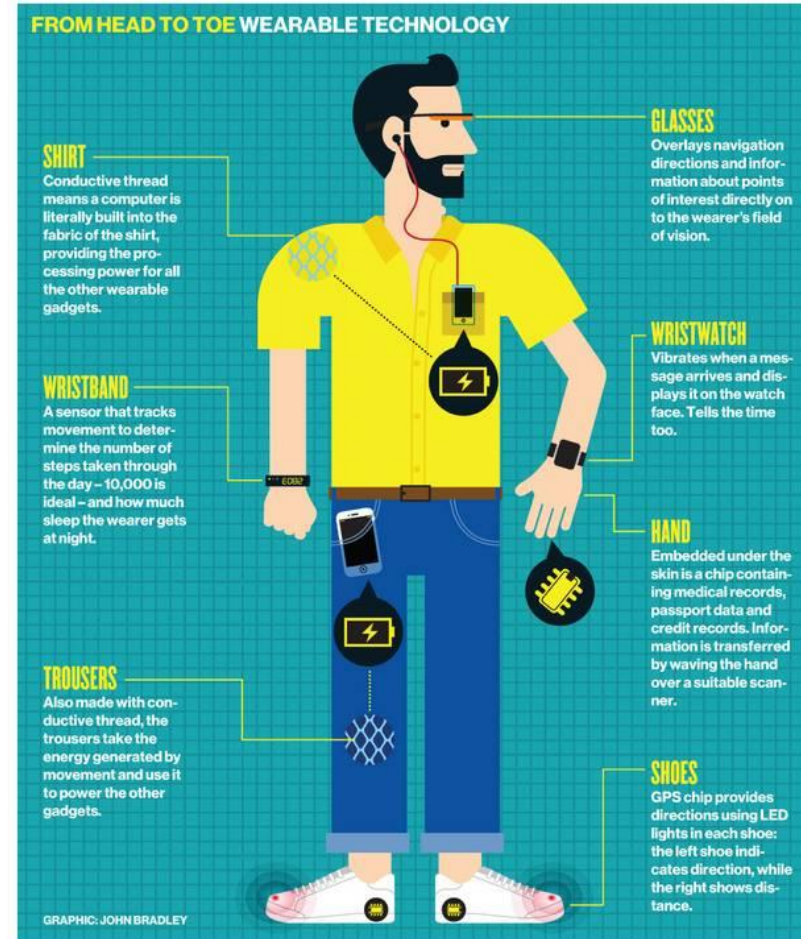
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Why do we need antennas? Motivation III



Galileo is Europe's own global navigation satellite

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❑ What is an antenna?

- All oscillating electric and magnetic fields propagate
- All circuit that create AC electric field and currents will radiate to some extent
- Antennas are devices to optimize and control the emitted radiation so you can couple electrical energy from a circuit to free space and back again

Antenna Design factor

- ❑ The strength of the radiated field in different direction (antenna pattern)
- ❑ The total power radiated compared to the driving power
- ❑ The impedance of the antenna to match it to transmission lines

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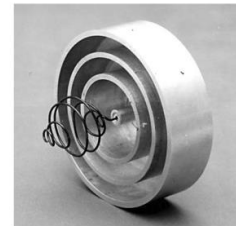
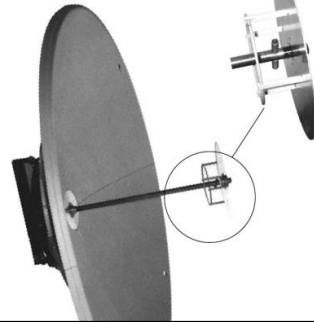
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- ❑ Potentially any conducting or dielectric structure can serve as an antenna...
- ❑ An antenna is designed to radiate or receive electromagnetic energy
But it is necessary
 - To design the antenna with **directional and polarization** properties suitable for the intended **application**
 - To **minimize energy reflection** in the transmission line-antenna juncture (to match the antenna impedance to the line Z_0)
 - The appropriate properties for the antenna are governed by its **shape and size** and the **material** of which it is made

❑ Key design factor

- Radiation Pattern
- Directivity
- Efficiency/ Gain
- Bandwidht



Spiral wire antenna on corrugated disk



Waveguide slot antenna array

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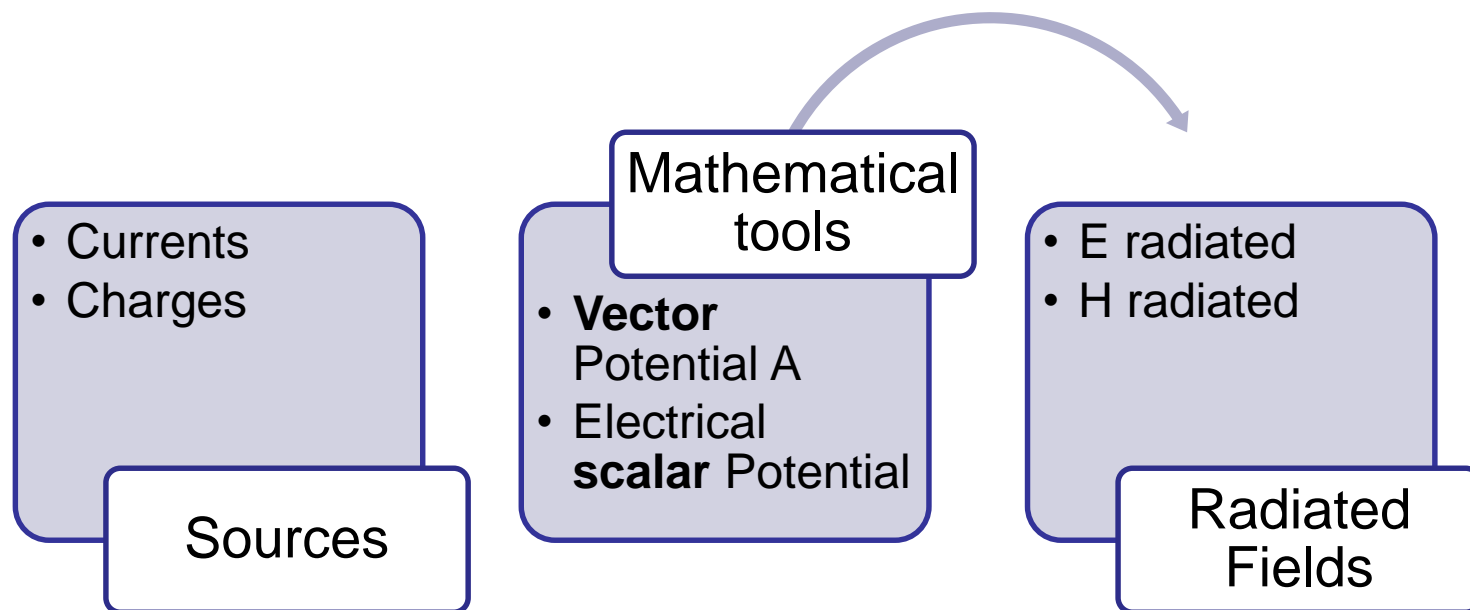
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❑ How we can obtain the radiated field?

- Not easy to solve (no simple wave equation)
- It is necessary to take into account the sources

❑ First solution of Maxwell equations with sources:

Radiated waves



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The vector potential \vec{A} for an electric current source \vec{J}

We start from the Maxwell equation $\nabla \cdot \vec{B} = 0$

Therefore, it can be represented as the curl of another vector $\nabla \cdot \nabla \times \vec{A} = 0$

$$\vec{B} = \nabla \times \vec{A} \quad \text{or} \quad \vec{H}_A = \frac{1}{\mu} \nabla \times \vec{A}$$

\vec{A} indicates the field due to \vec{A} potential

Substituting into Faraday's equation $\nabla \times \vec{E}_A = -j\omega\mu\vec{H}_A$

$$\left. \begin{aligned} \nabla \times [\vec{E}_A + j\omega\vec{A}] &= 0 \\ \nabla \times [-\nabla\Phi] &= 0 \end{aligned} \right\} \vec{E}_A + j\omega\vec{A} = -\nabla\Phi$$

Scalar function that represents an arbitrary **electric scalar potential**



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How \mathbf{A} is related to the sources?

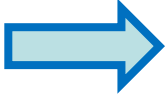
$$\begin{aligned}\vec{E}_A &= -\nabla\Phi - j\omega\vec{A} \\ \vec{B} &= \nabla \times \vec{A}\end{aligned}$$

$$\nabla \times \vec{H}_A = \vec{J} + j\omega\vec{E}_A$$

By replacing E and H in Maxwell equations by the former expressions,

...and using $\nabla \times \nabla \times \vec{A} = \nabla(\nabla \cdot \vec{A}) - \nabla^2 \vec{A}$

we obtain:


$$\nabla^2 \vec{A} + \omega^2 \mu \epsilon \vec{A} = -\mu \vec{J} + \nabla \left(\mu \epsilon j \omega \Phi + \nabla \cdot \vec{A} \right)$$

We need now to impose a condition to have a particular solution

$$j\omega \mu \epsilon \Phi + \nabla \cdot \vec{A} = 0$$

LORENTZ's condition or radiation condition

Finally

$$\nabla^2 \vec{A} + k^2 \vec{A} = -\mu \vec{J}$$

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How ϕ is related to the sources?

In the same way we can obtain for the potential ϕ

$$\nabla^2 \Phi + k^2 \Phi = -\frac{\rho}{\epsilon}$$

$$\nabla^2 \vec{A} + k^2 \vec{A} = -\mu \vec{J}$$

How are the solutions for these differential equations?

$$\vec{J} = J_z \quad \nabla^2 A_z + k^2 A_z = 0$$

$$\frac{d^2 A_z(r)}{dr^2} + \frac{2}{r} \frac{dA_z(r)}{dr} + k^2 A_z(r) = 0$$

homogeneous potential
Wave equation

Since in the limit the source is a point, A_z is not a function of direction (θ and ϕ)

Has two independent solutions

$$A_{z1} = C_1 \frac{e^{-jkr}}{r}$$

traveling waves

solutions

$$A_{z2} = C_2 \frac{e^{+jkr}}{r}$$

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Relation between sources and potentials

Electric current with harmonic time variation...

with sources... $\nabla^2 \vec{A} + k^2 \vec{A} = -\mu \vec{J}$

inhomogeneous Wave equation

$$\vec{J} = J_z \quad \nabla^2 A_z + k^2 A_z = -\mu J_z$$

$$A_z = \frac{\mu}{4\pi} \int_V J_z \frac{e^{-jkr}}{r} dv'$$

The sum of contribution of incremental \mathbf{J}_z sources to the potential

$$\left. \begin{aligned} \nabla^2 A_x + k^2 A_x &= -\mu J_x \\ \nabla^2 A_y + k^2 A_y &= -\mu J_y \\ \nabla^2 A_z + k^2 A_z &= -\mu J_z \end{aligned} \right\} \nabla^2 \vec{A} + k^2 \vec{A} = -\mu \vec{J}$$



we assume the solution

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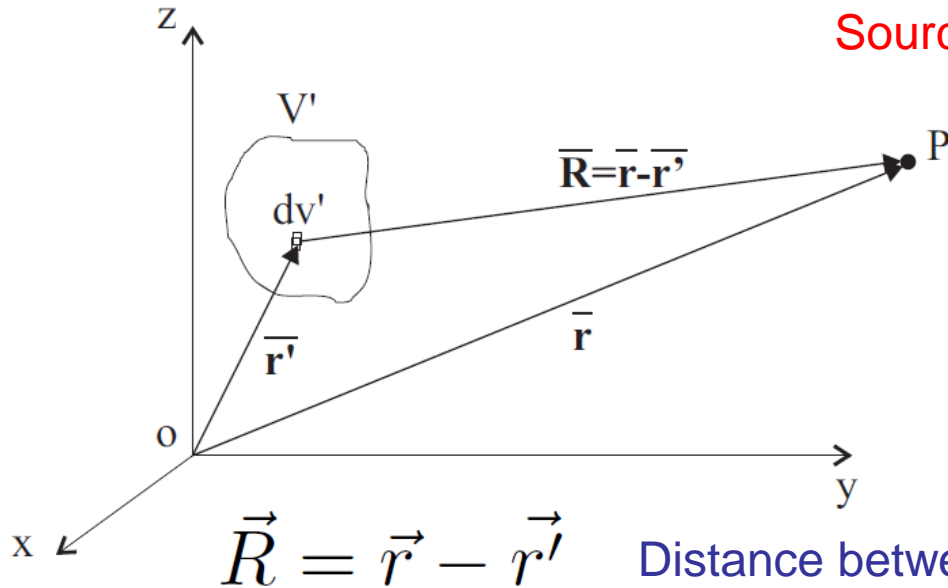
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ELECTROMAGNETIC RADIATION

Source not at origin, located in a volume V'



\vec{J} Electric current source

\vec{r}' Position of the source from the origin

\vec{r} Position of point P from the origin

$$\vec{A} = \frac{\mu}{4\pi} \int_{V'} \vec{J}(\vec{r}') \frac{e^{-jkR}}{R} dv'$$

We must integrate....

The contribution of the J with their directions and amplitudes and the...

e^{-jkR}

1

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In ...SUMMARY

we obtain the radiated fields in 3 steps....

1 Specify the sources J \vec{J}

2 Find vector potential A

$$\vec{A} = \frac{\mu}{4\pi} \int_{V'} \vec{J}(\vec{r}') \frac{e^{-jkR}}{R} dv'$$

3 Find the fields from the potential

$$\vec{E}_A = -j\omega\vec{A} - j \frac{\nabla(\nabla \cdot \vec{A})}{\omega\mu\epsilon}$$

$$\vec{H} = \frac{1}{\mu} \nabla \times \vec{A}$$

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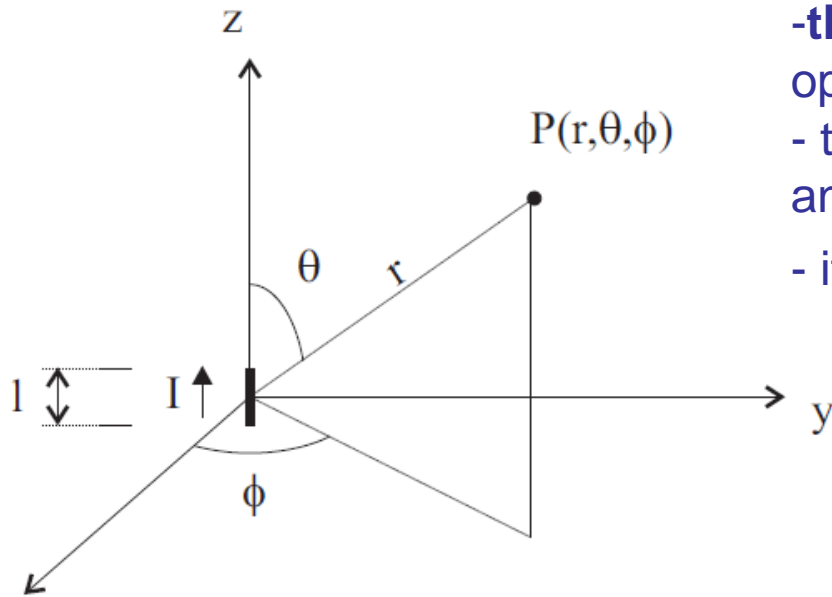
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CALCULATION OF THE RADIATED FIELDS FROM A SHORT DIPOLE

The Ideal dipole or Hertzian dipole



- the dipole** is very short compared to the operating wavelength
- the element of current is along the z-axis and centered on the coordinate origin
- it is of constant amplitude I_0

$$\vec{A} = \frac{\mu}{4\pi} \int_{V'} \vec{J} \frac{e^{-jkR}}{R} dv' = \frac{\mu}{4\pi} \int_C \vec{I} \frac{e^{-jkR}}{R} dl' = \hat{z} \frac{\mu I_0}{4\pi r} \frac{e^{-jkr}}{r} \int_{-l/2}^{l/2} dz' = \hat{z} \frac{\mu I_0 l}{4\pi r} e^{-jkr}$$

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Transformations Between Coordinate Systems

$$\begin{aligned}x &= r \sin \theta \cos \phi \\y &= r \sin \theta \sin \phi \\z &= r \cos \theta\end{aligned}$$

$$\begin{aligned}\hat{r} &= \hat{x} \cos \phi \sin \theta + \hat{y} \sin \phi \sin \theta + \hat{z} \cos \theta \\ \hat{\theta} &= \hat{x} \left[\cos \phi \cos \theta \right] + \hat{y} \left[\sin \phi \cos \theta \right] - \hat{z} \left[\sin \theta \right] \\ \hat{\phi} &= -\hat{x} \sin \phi + \hat{y} \cos \phi\end{aligned}$$

Unit vectors

$$\begin{aligned}\hat{x} &= \hat{r} \sin \theta \cos \phi + \hat{\theta} \cos \theta \cos \phi - \hat{\phi} \sin \phi \\ \hat{y} &= \hat{r} \sin \theta \sin \phi + \hat{\theta} \cos \theta \sin \phi + \hat{\phi} \cos \phi \\ \hat{z} &= \hat{r} \cos \theta - \hat{\theta} \sin \theta\end{aligned}$$

For example, to express the spherical components A_θ, A_ϕ in terms of the cartesian components, we proceed as follows:

$$A_\theta = \hat{\theta} \cdot \mathbf{A} = \hat{\theta} \cdot (\hat{x} A_x + \hat{y} A_y + \hat{z} A_z) = (\hat{\theta} \cdot \hat{x}) A_x + (\hat{\theta} \cdot \hat{y}) A_y + (\hat{\theta} \cdot \hat{z}) A_z$$

$$A_\phi = \hat{\phi} \cdot \mathbf{A} = \hat{\phi} \cdot (\hat{x} A_x + \hat{y} A_y + \hat{z} A_z) = (\hat{\phi} \cdot \hat{x}) A_x + (\hat{\phi} \cdot \hat{y}) A_y + (\hat{\phi} \cdot \hat{z}) A_z$$

The dot products can be read off

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RADIATED FIELDS FOR THE INFINITESIMAL DIPOLE (II)

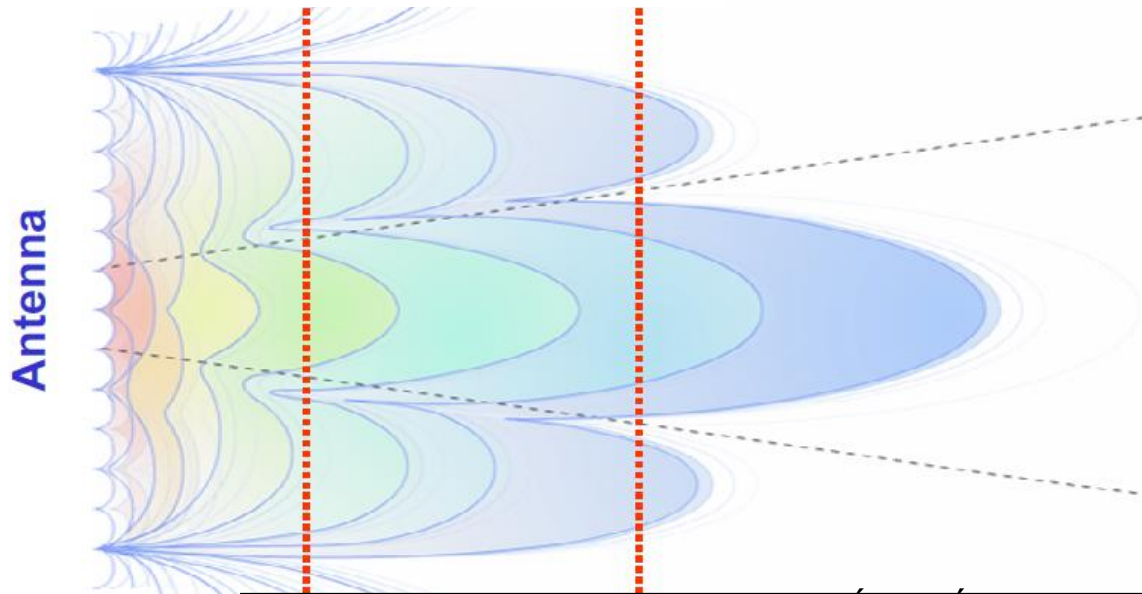
$$E_r = \eta \frac{I_0 l \cos \theta}{2\pi r^2} \left[1 + \frac{1}{jkr} \right] e^{-jkr}$$

$$E_\theta = j\eta \frac{k I_0 l \sin \theta}{4\pi r} \left[1 + \frac{1}{jkr} - \frac{1}{(kr)^2} \right] e^{-jkr}$$

$$E_\phi = 0$$

$$H_r = H_\theta = 0$$

$$H_\phi = j \frac{k I_0 l \sin \theta}{4\pi r} \left[1 + \frac{1}{jkr} \right] e^{-jkr}$$



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RADIATED FIELDS FOR THE INFINITESIMAL DIPOLE (III)

$$r \ll \lambda$$

We could neglect the terms depending of $\frac{r}{\lambda}$

$$E_r = -j\eta \frac{I_0 l \cos \theta e^{-jkr}}{2\pi k r^3}$$

$$E_\theta = -j\eta \frac{I_0 l e^{-jkr} \sin \theta}{4\pi k r^3}$$

$$E_\phi = 0$$

Near Field

$$H_r = H_\theta = 0$$

$$H_\phi = \frac{I_0 l e^{-jkr} \sin \theta}{4\pi r^2}$$

If $r \gg \lambda$ the predominant terms are

$$\frac{1}{r}$$

(The other terms can be neglected)

$$E_r = 0$$

$$E_\theta = j\eta \frac{k I_0 l \sin \theta e^{-jkr}}{4\pi r}$$

Far Field zone

$$H_r = 0$$

$$H_\theta = 0$$

$$H_\phi = \frac{k I_0 l \sin \theta e^{-jkr}}{4\pi r}$$

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In the Far-Field

- ❑ The **E** and **H** field component **are perpendicular** to each other and **transverse to the radial direction of propagation**
- ❑ The **E** and **H** field component are in **time phase**
- ❑ The E and H field component are both varying **inversely with the distance** to the dipole
- ❑ The **spherical wave** can be locally considered as a **plane wave**

with
$$\vec{H} = \frac{\hat{r} \times \vec{E}}{\eta}$$

in practice we can assume

$$\frac{E_{\theta}}{H_{\phi}} \approx \eta$$

We can compute the Poynting vector for this field...

$$\vec{W} = \frac{1}{R^2} \text{Re}(\vec{E} \times \vec{H}^*) = \hat{r} \frac{1}{R^2} |E_{\theta}|^2 = \hat{r} \eta |k I_0 l|^2 \sin^2 \theta$$

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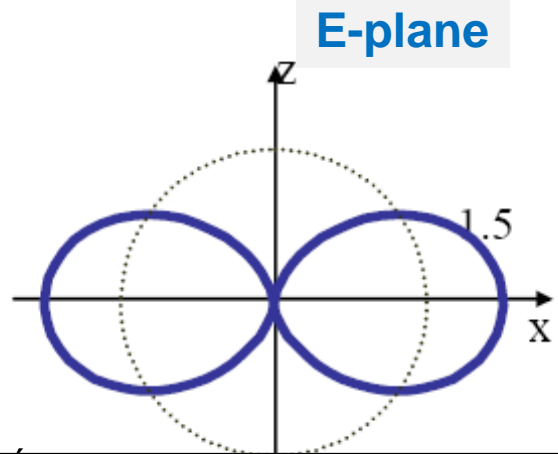
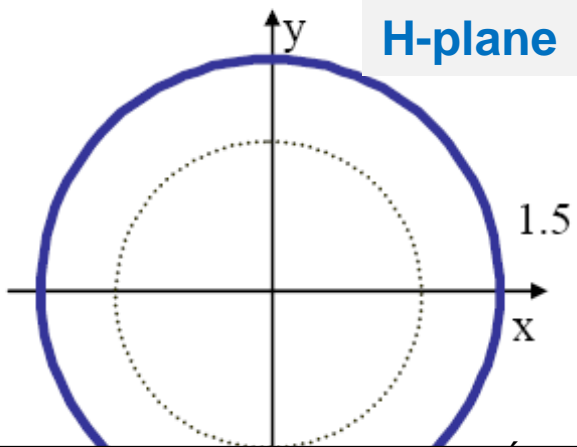
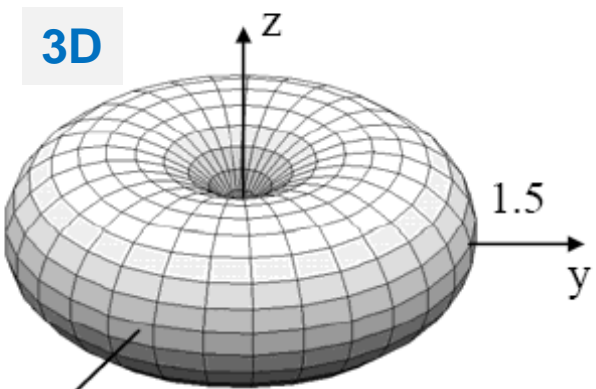
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□ Infinitesimal Dipole (short dipole)

□ How we can represented the radiated power?

- It's a function of spatial variables (θ, ϕ)
- 3D or 2D plots in specific planes
- No energy is radiated by the dipole along the direction of the dipole axis
- Maximum radiation occurs in $\theta=90^\circ$ (broadside direction)



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